

# Hydrogenated Oxygen-Deficient Blue Anatase as Anode for High-Performance Lithium Batteries

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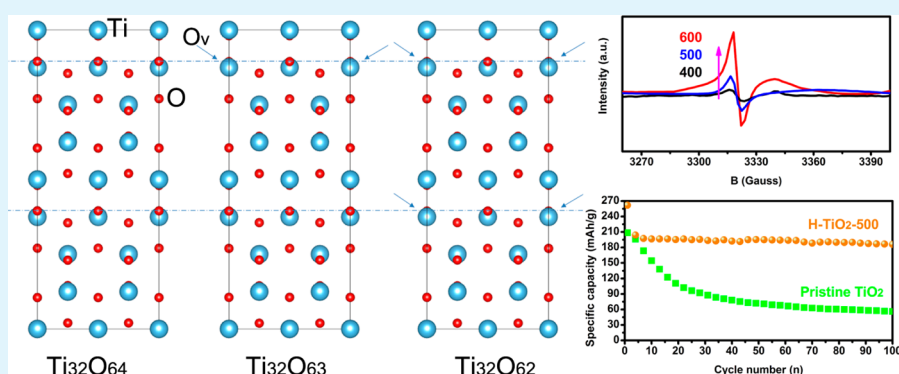
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## S Supporting Information



**ABSTRACT:** Blue oxygen-deficient nanoparticles of anatase  $\text{TiO}_2$  ( $\text{H-TiO}_2$ ) are synthesized using a modified hydrogenation process. Scanning electron microscope and transmission electron microscope images clearly demonstrate the evident change of the  $\text{TiO}_2$  morphology, from 60 nm rectangular nanosheets to much smaller round or oval nanoparticles of  $\sim 17$  nm, after this hydrogenation treatment. Importantly, electron paramagnetic resonance and positronium annihilation lifetime spectroscopy confirm that plentiful oxygen vacancies accompanied by  $\text{Ti}^{3+}$  are created in the hydrogenated samples with a controllable concentration by altering hydrogenation temperature. Experiments and theory calculations demonstrate that the well-balanced  $\text{Li}^+/\text{e}^-$  transportation from a synergetic effect between  $\text{Ti}^{3+}$ /oxygen vacancy and reduced size promises the optimal  $\text{H-TiO}_2$  sample a high specific capacity, as well as greatly enhanced cycling stability and rate performance in comparison with the other  $\text{TiO}_2$ .

**KEYWORDS:** anatase  $\text{TiO}_2$ , hydrogenation treatment, oxygen vacancy dependence, Li-storage performance, synergetic effect

## INTRODUCTION

Recently, titanium dioxide ( $\text{TiO}_2$ ) has attracted great attention as anode materials in rechargeable lithium-ion batteries (LIBs) due to its potential advantages, such as low cost, nontoxicity, high stability, and excellent cyclability. However, the low lithium-ion mobility and the poor electronic conductivity prevent  $\text{TiO}_2$  from practical application in LIBs.<sup>1,2</sup> Therefore, the development of simple and effective synthesis strategies for high-performance  $\text{TiO}_2$  anode materials is much desired.

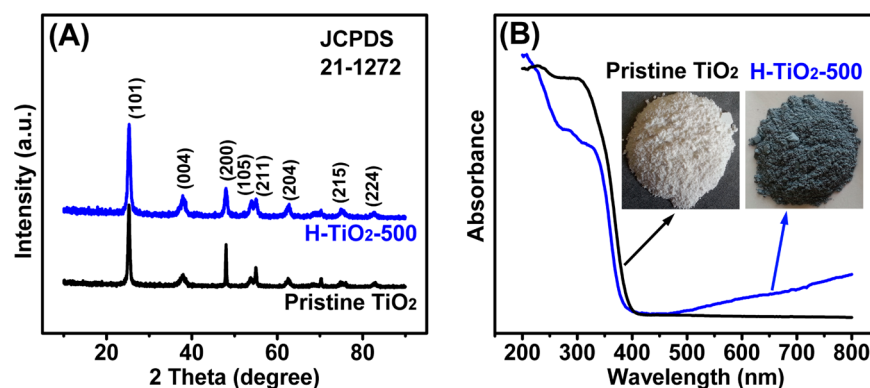
Design of nanostructured  $\text{TiO}_2$  has been proved effective in shortening the  $\text{Li}^+$  diffusion pathway, and the nanostructure (nanoparticles, nanowires, nanotubes, mesoporous materials, etc.) is critical to its electrochemical performance.<sup>3–7</sup> Coating or linking conductive components on/with  $\text{TiO}_2$ , such as  $\text{sp}^2$ -hybridized carbon, is another common route to modify the surface electronic properties of  $\text{TiO}_2$ , which is complex and

ineffective for large-scale production though.<sup>8–11</sup> People also establish techniques of doping  $\text{TiO}_2$  with foreign elements, such as niobium,<sup>12</sup> tin,<sup>13</sup> and nitrogen,<sup>14</sup> to regulate the band gap for enhanced electronic conductivity. However, the conductivity is still hardly sufficient for practical application. Recently, hydrogenation treatment has been developed as a facile technique to create plentiful oxygen vacancies in  $\text{TiO}_2$ , which can remarkably boost its electronic conductivity.<sup>15–17</sup> Rather than only producing a thin conductive layer on the surface, this  $\text{H}_2$  reduction method is able to turn the whole  $\text{TiO}_2$  electrode to highly conductive semiconductor.<sup>18</sup> In addition, this novel strategy is simple and can directly remove

Received: May 5, 2015

Accepted: October 12, 2015

Published: October 12, 2015



**Figure 1.** (A) X-ray diffraction patterns and (B) UV-vis reflectance spectra of the pristine TiO<sub>2</sub> and H-TiO<sub>2</sub>-500.

the oxygen atoms from TiO<sub>2</sub> to generate oxygen vacancies, which avoids inactive Li-storage constituents that may hinder the Li<sup>+</sup> diffusion. However, these oxygen vacancies in TiO<sub>2</sub> generate a large number of defects, which in turn deplete the concentration of Li<sup>+</sup> charge carriers, thus slowing the Li<sup>+</sup> diffusion. Therefore, to control the concentration of the oxygen vacancies is the key factor for fast Li<sup>+</sup> diffusion as well as high electronic conductivity.<sup>19,20</sup>

Hydrogenation treatment conducted at different conditions, such as atmospheric hydrogen mixture (5% H<sub>2</sub>+95% Ar, 1 atm, 450 °C), high-pressure pure hydrogen (H<sub>2</sub>, 20 bar, 200 °C or H<sub>2</sub>, 40 bar, 450 °C), etc., have been reported,<sup>21,22</sup> and the resulting materials exhibit various electrochemical properties. In this paper, we report a modified and relatively moderate hydrogenation process (pure H<sub>2</sub> under ca. 1 atm instead of under high-pressure environment), by which blue hydrogenated round anatase nanoparticles are obtained using rectangular anatase TiO<sub>2</sub> nanosheets as start materials. They are then used and tested as anode materials for LIBs. The effect of hydrogenation temperature on the structure defects and thus the electrochemical performance of the H-TiO<sub>2</sub> are investigated in detail. The electron paramagnetic resonance (EPR) results and the positron annihilation lifetime spectra (PALS) demonstrate that abundant oxygen vacancies and Ti<sup>3+</sup> are successfully introduced by this hydrogenation process. The blue anatase TiO<sub>2</sub> hydrogenated at 500 °C possesses significantly higher energy capacity as well as better cyclability and rate performance than the white pristine TiO<sub>2</sub>.

## EXPERIMENTAL SECTION

**Synthesis of Anatase TiO<sub>2</sub> Nanosheets.** A modified hydrothermal route was employed to synthesize anatase TiO<sub>2</sub> nanosheets.<sup>23</sup> In a typical experiment, 25 mL of Ti(OBu)<sub>4</sub> (Sinopharm Chemical Reagent Co., Ltd.) and 3 mL of HF (35.35 wt% Nanjing Chemical Reagent Co., Ltd.) were mixed in a 100 mL dried Teflon-lined autoclave and maintained at 200 °C for 24 h. After being cooled to room temperature, TiO<sub>2</sub> nanosheets with average length of 60 nm were obtained by high-speed centrifugation and then were washed with distilled water and absolute ethanol three times. The as-obtained TiO<sub>2</sub> is referred to as pristine TiO<sub>2</sub>.

**Hydrogenation Treatment of TiO<sub>2</sub> Nanoparticles.** Hydrogenation treatment was performed in an electric tube furnace equipped with a quartz tube reactor. The as-prepared TiO<sub>2</sub> nanosheet (1.2 g) was put into the tube furnace under a flow of pure Ar (purity 99.999%, 100 sccm) for 30 min to remove air in the furnace. Subsequently, the as-pretreated sample was heated to the set temperature (400, 500, and 600 °C) under pure hydrogen (99.999%) at a rate of 5 °C/min. The flow rate was fixed to 50 sccm, and the temperature was maintained for 0.5 h before it was

cooled to room temperature. These samples were referred to as H-TiO<sub>2</sub>-T (T = hydrogenation temperature [°C]). In addition, an oxygen-vacancy-free TiO<sub>2</sub> sample (denoted as TiO<sub>2</sub>-17 nm-R, R means remove the oxygen vacancy) with the same size as H-TiO<sub>2</sub>-500 was also synthesized via second heat treatment of H-TiO<sub>2</sub>-500 at 500 °C for 10 h under air to remove oxygen vacancy but keep its original structure.

**Computational Method.** In our theoretical calculations, TiO<sub>2</sub> lattice was modeled using a 2 × 2 × 2 supercell with a molecular formula of Ti<sub>32</sub>O<sub>64</sub>. Calculations were performed using the Vienna Ab-initio Simulation Package (VASP),<sup>24</sup> which employed a density functional theory (DFT)-based plane-wave method.<sup>25</sup> Exchange correlation effects were included using the Perdew–Burke–Ernzerhof (PBE).<sup>26</sup> The energy cutoff for the expansion of wave function was set to be 380 eV, which was tested to be enough for both total energy in the systems. For geometry relaxation, 4 × 4 × 2 Monkhorst–Pack K-points were adopted for the calculation. To correct the on-site electron correlation, DFT plus Hubbard model (DFT+U) was employed, with U = 4.0 eV based on our tests and early publications.<sup>27</sup> The atomic position was optimized through PBE calculations with a criterion that required the calculated forces less than 1 × 10<sup>-2</sup> eV/Å.

**Structural Characterization.** The anatase TiO<sub>2</sub> nanosheets before and after hydrogenation treatment were extensively characterized by X-ray diffraction analysis (XRD, Bruker D8 Advance X-ray diffractometer using a Cu Kα X-ray source), field emission scanning electron microscopy (FESEM, HITACHI S4800), high-resolution transmission electron microscopy (HRTEM, JEOL JEM 2100F), X-ray photoelectron spectrometry (XPS, Thermo ESCALAB 250), Raman spectroscopy (Raman, JY HR800), UV-vis diffuse reflectance spectra (UV-vis, Shimadzu UV = 3600), electron paramagnetic resonance (EPR, Bruker EMX-10/12), and positron annihilation technique spectroscopy (PALS, slow positron beam at institute of high energy physics, resolution = 197 ps).

**Electrochemical Measurement.** To prepare working electrode, the as-prepared pristine TiO<sub>2</sub> or H-TiO<sub>2</sub> materials, conductive carbon black, and polyvinylidene difluoride (PVDF) were mixed at a weight ratio of 75:15:10 in appropriate amount of N-methylpyrrolidinone (NMP) to obtain homogeneous yogurtlike slurry. The slurry was coated onto a thin copper foil, followed by drying at 80 °C for 12 h under vacuum. Electrode discs (diameter = 12 mm) were punched from the Cu foil, on which the typical loading of active material was ~1.0 mg/cm<sup>2</sup>. CR2032 coin cells were assembled in an Ar-filled glovebox using pure lithium metal foils as counter electrode, a polypropylene microporous film as separator, 1 M LiPF<sub>6</sub> dissolved in a mixture of ethylene carbonate (EC), and dimethyl carbonate (DMC) 1:1 (w/w) as electrolyte. Galvanostatic charge/discharge tests were performed with LAND CT-2001A instrument (Wuhan, PRC) at a voltage range of 1.0–3.0 V versus Li/Li<sup>+</sup> at room temperature (23 ± 0.5 °C). Cyclic voltammetry (CV) and electrochemical impedance spectroscopy (EIS) were performed on a CHI660D electrochemical workstation (CH Instrument, Shanghai, PRC).

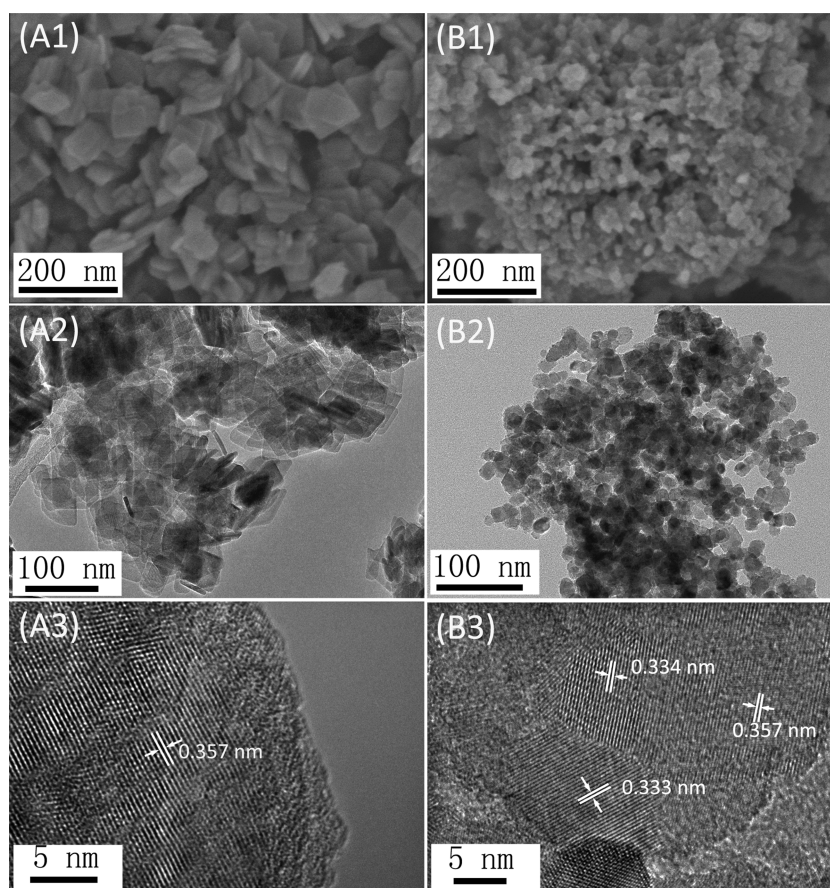


Figure 2. FESEM (A1, B1), TEM (A2, B2), and HRTEM (A3, B3) images of the (A) pristine  $\text{TiO}_2$  and (B)  $\text{H-TiO}_2\text{-500}$ .

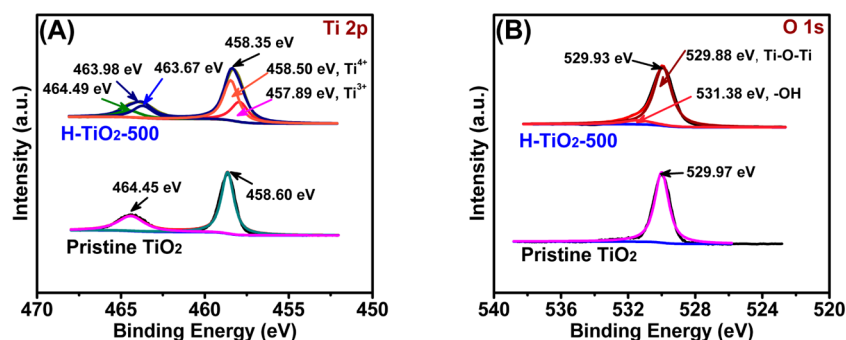


Figure 3. High-resolution XPS spectra of Ti 2p (A) and O 1s (B) of the pristine  $\text{TiO}_2$  and  $\text{H-TiO}_2\text{-500}$ .

## RESULTS AND DISCUSSION

Figure 1A shows the XRD patterns of pristine  $\text{TiO}_2$  and  $\text{H-TiO}_2\text{-500}$ . Both samples exhibit typical diffraction peaks of anatase phase (JCPDS 21–1272) at  $25.28^\circ$ ,  $48.05^\circ$ ,  $55.06^\circ$ ,  $37.80^\circ$ ,  $53.89^\circ$ , and  $62.69^\circ$ , corresponding to the planes of (101), (200), (211), (004), (105), and (204), respectively. No other phases such as rutile  $\text{TiO}_2$  or other reduced phases are observed, suggesting no obvious changes of the anatase framework after hydrogenation treatment. However, the color of  $\text{TiO}_2$  is visually translated from white to blue, indicating that microstructure is changed to some extent, which further affects its optical properties. The UV–vis diffraction spectra of pristine  $\text{TiO}_2$  and  $\text{H-TiO}_2\text{-500}$  in Figure 1B show that they have similar signal response in UV region (200–400 nm), but a much larger absorption is observed on  $\text{H-TiO}_2\text{-500}$  in the visible light range (400–800 nm). The optical changes of anatase  $\text{TiO}_2$  after

hydrogenation treatment could be attributed to the local  $\text{Ti}^{3+}$  centers (band at 620 nm) and oxygen vacancies (bands at 1060, 441, and 486 nm).<sup>5,28</sup>

The SEM and TEM images of pristine  $\text{TiO}_2$  (Figure 2A1–A3) present uniform rectangular nanosheets with an average side length of  $\sim 60$  nm. After hydrogenation treatment, it can be obviously observed that these nanosheets were turned into round or oval nanoparticles with an average size of  $\sim 17$  nm (Figure 2B1–B3 and Figure S1). It indicates that the original thin nanosheet structure of  $\text{TiO}_2$  is unstable under high-temperature hydrogen atmosphere, thus collapsing and reconstructing during the hydrogenation treatment. The resultant nanoparticles consist of many small grains with abundant grain boundaries. Furthermore, in addition to typical lattice fringes of (101) plane of anatase  $\text{TiO}_2$  with a lattice distance of 0.357 nm, a new lattice fringe with a lattice distance



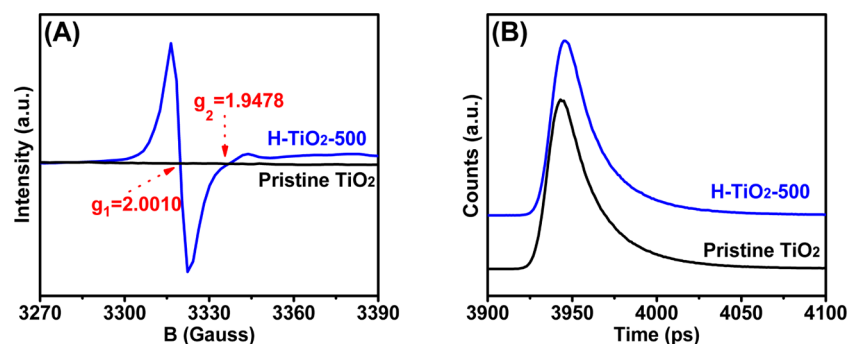


Figure 4. Low-temperature EPR (A) and PALS (B) spectra of pristine TiO<sub>2</sub> and H-TiO<sub>2</sub>-500.

Table 1. Positron Lifetimes and Relative Intensities with Their Deviations for Pristine TiO<sub>2</sub> and H-TiO<sub>2</sub>-T Samples at Different Temperature<sup>a</sup>

samples		$\tau_1$ (ps)	$\tau_2$ (ps)	$\tau_3$ (ns)	$I_1$ (%)	$I_2$ (%)	$I_3$ (%)	$I_1/I_2$
pristine TiO <sub>2</sub>	value	175.5	357.0	2.5820	46.10	49.30	4.66	0.94
	error	3.6	5.1	0.0240	1.40	1.40	0.14	
H-TiO <sub>2</sub> -400	value	191.2	398.0	2.5950	60.80	34.80	4.40	1.75
	error	3.1	7.9	0.0270	1.30	1.30	0.13	
H-TiO <sub>2</sub> -500	value	182.4	385.8	2.5740	63.30	32.00	4.69	1.98
	error	2.7	8.2	0.0330	1.30	1.30	0.13	
H-TiO <sub>2</sub> -600	value	173.6	356.5	2.6210	52.00	43.50	4.43	1.19
	error	3.6	6.0	0.0320	1.40	1.40	0.14	

<sup>a</sup>T = 400, 500, 600 °C.

of 0.330 nm appears in H-TiO<sub>2</sub>-500, corresponding to the (004) planes of Ti<sub>6</sub>O<sub>11</sub>. These results suggest that the hydrogenation process has made obvious deviation from the stoichiometry for TiO<sub>2</sub>, possibly leading to the introduction of oxygen vacancies into TiO<sub>2</sub>.

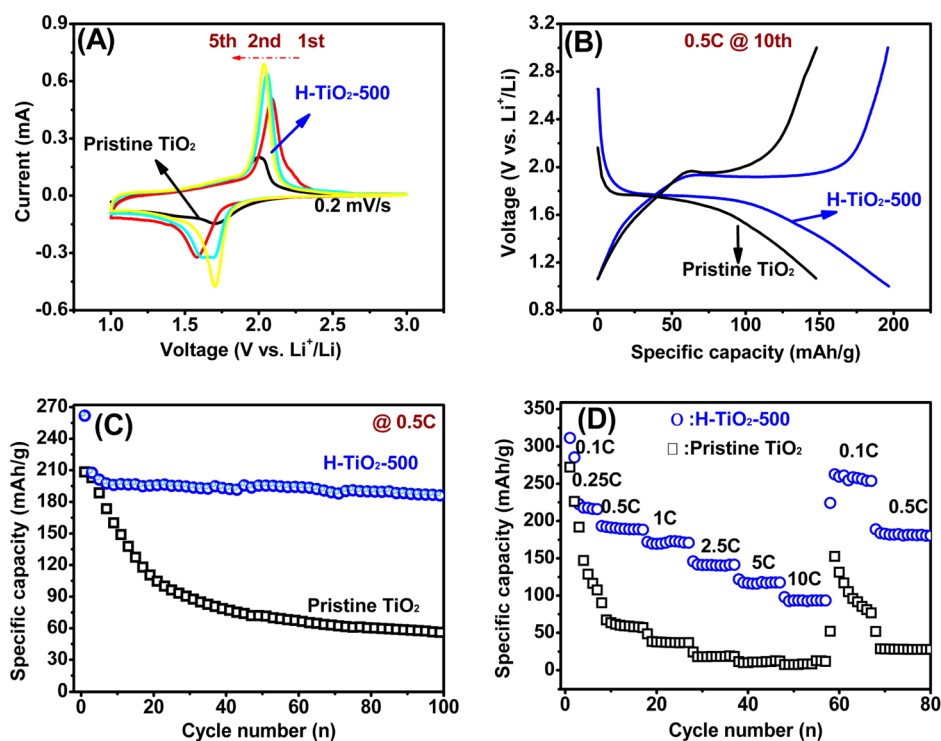
To deeply reveal the oxygen vacancy, multiple techniques, including XPS, EPR, and PALS, were employed to characterize the as-prepared samples. The original high-resolution Ti 2p XPS spectra in Figure S2 demonstrate that the full width at half-maximum (fwhm) for H-TiO<sub>2</sub>-500 and pristine TiO<sub>2</sub> is 1.475 and 0.915 eV, respectively. Furthermore, a small shift of binding energy from 458.60 eV of pristine TiO<sub>2</sub> to 458.35 eV of H-TiO<sub>2</sub>-500 is observed. Because these data are recorded on the same instrument under the same condition, we believe that these changes are probably attributed to the different chemical bonding environment of Ti 2p.<sup>29,30</sup> As shown in Figure 3A, after well deconvolution, a typical Ti<sup>4+</sup> 2p<sub>3/2</sub> signal peak at 458.60 and 458.50 eV for pristine TiO<sub>2</sub> and H-TiO<sub>2</sub>-500 is observed, respectively. Importantly, an extra obvious signal at 457.89 eV is observed in H-TiO<sub>2</sub>-500, corresponding to Ti<sup>3+</sup> 2p<sub>3/2</sub>. In contrast, there is no obvious Ti<sup>3+</sup> species observed in Ti 2p XPS spectrum of the pristine TiO<sub>2</sub>. Similar results are found on the high-resolution O 1s XPS spectra shown in Figure 3B. Besides the similar signal peak of Ti–O–Ti bonds located at ~529.9 eV for both samples, an extra peak centered at ~531.38 eV is observed for H-TiO<sub>2</sub>-500, which is attributed to Ti–OH species on the surface.<sup>31</sup> To some extent, these results of XPS analysis probably confirm that H<sub>2</sub> treatment has partially reduced Ti<sup>4+</sup> to Ti<sup>3+</sup> and created oxygen deficiencies.

To further identify the nature of radical species and the presence of oxygen vacancy/Ti<sup>3+</sup> ions, low-temperature EPR spectra of the pristine TiO<sub>2</sub> and H-TiO<sub>2</sub>-500 are recorded, as shown in Figure 4A. It has been reported that oxygen vacancy has a typical *g*-value of ~2.004 and paramagnetic Ti<sup>3+</sup> at *g* = 1.94–1.99.<sup>32</sup> Here, only the H<sub>2</sub>-treated sample gives rise to very

strong EPR signals, which can be fitted by two components with *g* = 2.0010 and 1.9478, corresponding to oxygen vacancy and Ti<sup>3+</sup>, respectively. Moreover, the room-temperature EPR spectrum of H-TiO<sub>2</sub>-500 shows smooth signals in the range of *g* = 1.94–1.99 of paramagnetic Ti<sup>3+</sup>, as Ti<sup>3+</sup> ions cannot be detected by EPR at room-temperature (Figure S3). The above results provide an effective proof for the existence of oxygen vacancy as well as Ti<sup>3+</sup>.

Positron annihilation lifetime spectrum (PALS), which is useful at inquiring and determining vacancy-type defects in semiconductors, is employed here to provide information on the positron lifetimes and thus the concentration of vacancies.<sup>33,34</sup> Figure 4B illustrates typical positron PALS spectra of pristine TiO<sub>2</sub> and H-TiO<sub>2</sub>-500. The positron lifetime is composed of three components, which can be evaluated by deconvolution of the positron decay spectrum *D*(*t*) using PATFIT program. The values of the three positron lifetime components, namely,  $\tau_1$ ,  $\tau_2$ , and  $\tau_3$ , and their corresponding intensities *I*<sub>1</sub>, *I*<sub>2</sub>, and *I*<sub>3</sub> are listed in Table 1. The first short lifetime of positrons  $\tau_1$  is generally due to the free annihilation of positrons in a defect-free crystal. However, some smaller vacancies and/or shallow positron traps in disordered crystals may reduce the surrounding electron density, leading to bigger values of  $\tau_1$ . This might account for the fact that the  $\tau_1$  of H-TiO<sub>2</sub>-500 (182.4 ps) is remarkably longer than that of pristine TiO<sub>2</sub> (175.5 ps). The longest lifetime component ( $\tau_3$ ) probably depends on the loading method of “sandwich” structure (i.e., the substrate material, the inevitable interface of specimen-positron and source-specimen), which has no useful information for revealing the defects of as-prepared TiO<sub>2</sub>.

The intermediate lifetime component  $\tau_2$  is the most important one and can be attributed to positrons captured by the low electron density area or defects. The average electron density in larger-sized defects is lower than that in smaller defects, which decreases the annihilation rate and increases the



**Figure 5.** Electrochemical tests of pristine TiO<sub>2</sub> and H-TiO<sub>2</sub>-500. (A) CVs at a scan rate of 0.2 mV/s (voltage range of 1.0–3.0 V); (B) galvanostatic charge/discharge voltage profiles at 0.5 C between voltage limits of 1 and 3 V (data taken from the 10th cycle); (C) cycling performance up to 100 cycles at 0.5 C; (D) specific discharge capacity at various C-rates of 0.1, 0.25, 0.5, 1, 2.5, 5, and 10 C. (1 C = 336 mA/g).

positron lifetime, correspondingly. Therefore, the values of  $\tau_2$  are much larger than those of  $\tau_1$ . Meanwhile, the ratio of  $I_1$  to  $I_2$  ( $I_1/I_2$ ) is used to compare the relative concentration of small-size bulk defects to large-size surface defects. It is also found that the  $\tau_2$  of H-TiO<sub>2</sub>-500 (385.8 ps) is much longer than that of pristine TiO<sub>2</sub> (357.0 ps), indicating reduced surrounding electron density in the former structure. Moreover, the ratio  $I_1/I_2$  of H-TiO<sub>2</sub>-500 is as high as 1.98 more than double of that of pristine TiO<sub>2</sub> ( $I_1/I_2 = 0.94$ ). Such results clearly demonstrate that some larger-size vacancy clusters are induced into the H-TiO<sub>2</sub>-500, which results from interaction between the small neutral Ti<sup>3+</sup>-oxygen vacancy associates.<sup>33</sup>

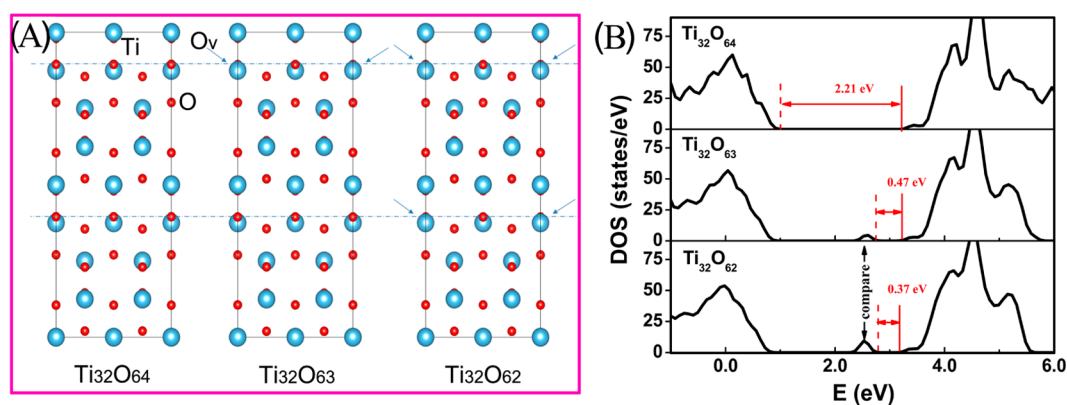
The pristine and H-TiO<sub>2</sub>-500 samples are used as anode materials to assemble half lithium batteries and undergo a series of electrochemical tests under identical conditions. The CVs (Figure 5A) show dramatic differences between them: a pair of pronounced redox peaks, situated at 1.571/2.095 V, is observed on the first cycle for H-TiO<sub>2</sub>-500, which is associated with the insertion/deinsertion processes of Li<sup>+</sup> in the TiO<sub>2</sub> framework, following the redox reaction:  $\text{TiO}_2 + x\text{Li}^+ + xe^- \leftrightarrow \text{Li}_x\text{TiO}_2$ , while a minor reduction and oxidation peak appears on the CV curves for pristine TiO<sub>2</sub>. Meanwhile, from the second cycle onward, as shown in Figure 5A, the CV curves for H-TiO<sub>2</sub>-500 are pretty reproducible, indicating a good reversibility for the Li<sup>+</sup> intercalation/deintercalation process during cycling.

The galvanostatic discharge/charge curves of the 10th cycle at 0.5 C for pristine TiO<sub>2</sub> and H-TiO<sub>2</sub>-500 are illustrated in Figure 5B. Both the TiO<sub>2</sub> materials display typical cathodic and anodic potential plateau characteristic of anatase TiO<sub>2</sub> at ~1.75 and 1.95 V versus Li/Li<sup>+</sup>. However, H-TiO<sub>2</sub>-500 has a greatly enhanced discharge specific capacity of 196.1 mAh/g and columbic efficiency of 99.7%, while pristine TiO<sub>2</sub> only achieved 147.5 mAh/g and 95.6% at 0.5 C. As shown in Figure 5C, H-

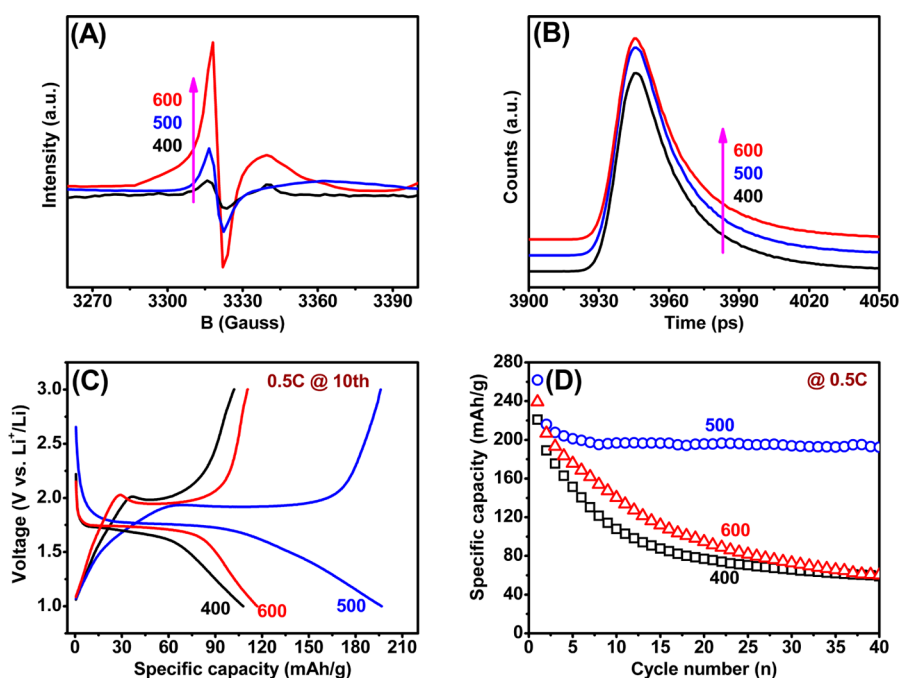
TiO<sub>2</sub>-500 demonstrates much more stable cycling performance, and the discharge capacity can be retained at 186 mAh/g after 100 cycles, while only 55.3 mAh/g can be kept after 100 cycles for pristine TiO<sub>2</sub>. The improved cycling performance of the blue H-TiO<sub>2</sub>-500 can be credited to the decrease in particle sizes, crystalline dislocation, and presence of oxygen vacancy/Ti<sup>3+</sup> induced by the hydrogenation process.

Figure 5D shows the rate capability of H-TiO<sub>2</sub>-500 and pristine TiO<sub>2</sub> with current densities varying from 0.1 C to 10 C. Each of the two samples exhibits high initial discharge capacity at 0.1 C; however, a sharp capacity decay occurs in the following two cycles for the pristine TiO<sub>2</sub>, while the H-TiO<sub>2</sub>-500 retains dramatically higher discharge capacity. As the increasing the C rate, the pristine TiO<sub>2</sub> undergoes even further decay of the capacities, but the H-TiO<sub>2</sub>-500 anode material can endure more radical changes of current density. Even after 60 cycles with successively increasing current densities from 0.1 to 10 C, H-TiO<sub>2</sub>-500 can still retain a discharge capacity of ~100 mAh/g, and when the rate returns to 0.1 C, the discharge capacity reverses to the original value of 275 mAh/g at 0.1 C. These results suggest that the advantages of hydrogenation treatment on TiO<sub>2</sub> in respect of the cyclability are more significant.

Moreover, it can be found that our prepared H-TiO<sub>2</sub>-500 shows some advantages in electrochemical capacity under different conditions over typical nanoporous anatase TiO<sub>2</sub> mesocrystals,<sup>7</sup> traditionally hydrogenated blue rutile TiO<sub>2</sub>,<sup>15</sup> and anatase TiO<sub>2- $\delta$</sub> .<sup>19</sup> (Figure S4). Such significant improvement on the electrochemical properties can be attributed to the enhanced electronic conductivity as well as Li<sup>+</sup> diffusion by hydrogenation treatment. As shown in Figure S11, the EIS results demonstrate that H-TiO<sub>2</sub>-500 has lowest electronic resistance (R<sub>ct</sub>) and Warburg resistance (W), which can be



**Figure 6.** Density functional theory theoretical calculations: (A) Calculated models of  $\text{Ti}_{32}\text{O}_{64}$  (bulk  $\text{TiO}_2$ ),  $\text{Ti}_{32}\text{O}_{63}$  (1.56% oxygen deficiency), and  $\text{Ti}_{32}\text{O}_{62}$  (3.13% oxygen deficiency) (arrow: oxygen vacancy); (B) the total density states for the three  $\text{TiO}_2$  models. The red dash lines mean the Fermi level in three difference model.



**Figure 7.** (A) EPR spectra, (B) PALS spectra, (C) galvanostatic charge/discharge profiles (data taken from the 10th cycle), (D) cycling performance up to 100 cycles at 0.5 C of H- $\text{TiO}_2$  prepared at different hydrogenation temperatures.

probably ascribed from abundant oxygen vacancies and  $\text{Ti}^{3+}$  created by hydrogenation treatment in  $\text{TiO}_2$  for improved electrical conductivity and the spitting distance between the oxygen vacancies and the  $\text{Ti}^{3+}$  cores that could repel the positively charged  $\text{Li}^+$  and deliver them into other  $\text{Ti}^{3+}$  cores to accelerate the  $\text{Li}^+$  transport.<sup>35</sup>

Note that the decreased particle size also plays an important role in shortening distance for  $\text{Li}^+$  transport and accelerate  $\text{Li}^+$  transport. To verify this point, we synthesize an oxygen-vacancy-free  $\text{TiO}_2$  (denoted as  $\text{TiO}_2$ -17 nm-R) with a same size as H- $\text{TiO}_2$ -500 via second heat treatment of H- $\text{TiO}_2$ -500 under air to remove oxygen vacancy but keep its structure (Figures S5–S7). The cycling performance tests (Figure S8) present that the performance of  $\text{TiO}_2$ -17 nm-R sample is little higher than pristine  $\text{TiO}_2$  but far from H- $\text{TiO}_2$ . Therefore, we can conclude that it is the synergistic effect from oxygen vacancy/ $\text{Ti}^{3+}$  and decreased size that enhances the electronic

conductivity as well as  $\text{Li}^+$  transport, resulting in a high-capacity and stable  $\text{TiO}_2$  material for Li-ion battery.

Furthermore, the mechanism of oxygen-deficiency/ $\text{Ti}^{3+}$  on the electron conductivity of  $\text{TiO}_2$  is studied by DFT calculations. As shown in Figure 6A, three calculation models of large supercell, including  $\text{Ti}_{32}\text{O}_{64}$  (bulk  $\text{TiO}_2$ ),  $\text{Ti}_{32}\text{O}_{63}$  (1.56% oxygen vacancy), and  $\text{Ti}_{32}\text{O}_{62}$  (3.13% oxygen vacancy), are established through removing lattice oxygen between two interstitial voids in the hydrogenation reduction process. The red balls stand for oxygen atom and green for titanium atom; the arrow indicates the missing location of O atom, namely, the location of oxygen-deficiency. Projection of the  $2 \times 2 \times 2$   $\text{TiO}_2$  supercell is on the ac-plane (the corresponding stereogram is shown in Figure S9). It can be found from the total density of states (DOSs) in Figure 6B that the Fermi level increases from  $\sim 1.00$  eV for  $\text{TiO}_2$  to  $\sim 2.76$  eV for  $\text{Ti}_{32}\text{O}_{63}$  and 2.79 eV to  $\text{Ti}_{32}\text{O}_{62}$ . Moreover, the calculated intrinsic band gap ( $E_g$ ) for pristine  $\text{TiO}_2$  is 2.21 eV, which is in line with previous

experimental and computational results,<sup>36,37</sup> while the calculated  $E_g$  values of  $Ti_{32}O_{63}$  and  $Ti_{32}O_{62}$  dramatically drop to 0.47 and 0.37 eV, respectively. Therefore, it can be concluded that the electronic conductivity can be enhanced via adjusting the properties of  $E_g$  by oxygen vacancies. In addition, comparing the DOSs of  $Ti_{32}O_{63}$  and  $Ti_{32}O_{62}$ , more DOSs are found below Fermi level in  $Ti_{32}O_{62}$  than that in  $Ti_{32}O_{63}$ , suggesting a higher charge carrier concentration with increase of oxygen vacancies. The above results illustrate that the decrease of band gap and the increase of DOSs below Fermi level in the presence of oxygen vacancies might be the two reasons for the improvement of electronic conductivity of the hydrogenated  $TiO_2$ .

The effect of hydrogenation temperature on electrochemical properties of  $H_2$ -treated  $TiO_2$  is also studied. A series of  $H_2$  treated anatase (XRD in Figure S10) are prepared at different hydrogenation temperatures of 400, 500, and 600 °C. Figure 7A is their low-temperature EPR spectra. The typical peaks at gauss values of ca. 3366 associated with the signals of oxygen vacancies are observed for all the three  $H_2$ -treated samples, and the concentration is gradually increased as the hydrogenation temperature is elevated from 400 to 600 °C. In addition, the positron lifetime component  $\tau_2$  reduces with increasing the hydrogenation temperature as compared in Table 1, and Figure 7B is the initial graph of matching PALS spectra. The EIS results in Figure S11 and Table S1 demonstrate that the electronic resistance ( $R_{ct}$ ) of  $TiO_2$  after hydrogenation at 400 or 600 °C is higher than that at 500 °C, suggesting that 500 °C is the optimal condition for the improvement of electronic conductivity. Furthermore, the Warburg resistance ( $W$ ) is obviously reduced for all hydrogenated  $TiO_2$  samples, indicating that the ionic diffusion process inside electrode is accelerated significantly probably due to the shortened Li-ion transport path resulting from the reduced size under hydrogenation process. The above results reveal that the oxygen vacancies mostly determine the electronic conductivity for hydrogenated  $TiO_2$ .

Interestingly, the electrochemical results in Figure 7C,D demonstrate that there is not an obvious trend of linear relationship between the Li-storage performance and the concentration of oxygen vacancies. It can be seen that the specific discharge capacity of H- $TiO_2$ -500 is much higher than that of H- $TiO_2$ -400 and H- $TiO_2$ -600, and the former also exhibits much better cycling performance than the latter two, as shown in Figure 7D. The relationship between the lithium storage properties and the concentration of the oxygen vacancies has been discussed by Shin et al.<sup>19</sup> who declared that well-balanced  $Li^+/e^-$  transport is the key factor for high-performance  $TiO_2$  anode. In their study, the prolonged hydrogenation treatment time results in a greater electronic conductivity but a smaller  $Li^+$  conductivity due to the higher concentration oxygen vacancies. A proper degree of reduction must be desired for high-performance electrode materials. Our results also reveal that in spite of more oxygen vacancies generated at higher hydrogenation temperatures; the  $Li^+$  diffusion is not obviously enhanced due to the depression of free  $Li^+$  concentration by excess electrons. As a consequence, the  $TiO_2$  hydrogenated at 500 °C in this situation has the well-balanced  $Li^+/e^-$  transport, resulting in a greatest improvement over Li-storage performance.

## CONCLUSION

In summary, blue oxygen vacancies anatase  $TiO_2$  anodes were synthesized using a modified hydrogenation process for LIBs. Systematical characterizations reveal that  $Ti^{3+}$ /oxygen vacancy was successfully injected in  $TiO_2$  with a controllable concentration by hydrogenation temperature. Further experiments and theory calculation demonstrate that oxygen vacancies can act on the band gap and then decrease the  $E_g$  and increase the DOSs below Fermi level, which facilitate the enhancement of the electronic conductivity of  $TiO_2$ . Meanwhile, the spitting distance between the oxygen vacancies and the  $Ti^{3+}$  cores could repel the positive charged  $Li^+$  and deliver them into other  $Ti^{3+}$  cores to accelerate the  $Li^+$  transport. Further accompanied with the synergetic effect from the reduced sized-derived shortened transport path distance, the H- $TiO_2$ -500 with a well-balanced  $Li^+/e^-$  transport possesses a high specific capacity, as well as greatly enhanced cycling stability and rate performance in comparison with the other  $TiO_2$  samples, endowing it as a potential high-performance anode for lithium batteries.

## ASSOCIATED CONTENT

### Supporting Information

The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acsami.5b07000.

Relevant XRD, TEM, particle size distributions, original Ti 2p XPS curves, low- and room-temperature EPR, EIS results and additional cycle performance for  $TiO_2$ , H- $TiO_2$  and/or  $TiO_2$ -17 nm-R; electrochemical property comparison, and DFT calculation mode. (PDF)

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### Author Contributions

J.Z., G.B.J., and Y.S.L. conceived and designed the study; J.Z. performed the experiments, and J.Z., Y.S.L., and G.B.J. analyzed the data; P.Z., X.Z.C., and B.Y.W. provided the PALS test; X.G.Z., Y.Z., and D.N.S. performed the density functional theory (DFT) calculations. J.Z. wrote the paper; J.Z., C.H.Z., and G.B.J. reviewed and edited the manuscript. All authors read and approved the manuscript.

### Notes

The authors declare no competing financial interest.

## ACKNOWLEDGMENTS

This work was financially supported by the National Natural Science Foundation of China (No. 11575085), National Basic Research Program of China (973 program, No: 2013CB934700), and the Priority Academic Program Development of Jiangsu Higher Education Institutions.

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